

Clouds and turbulent moist convection

Lecture 2: Cloud formation and Physics

Caroline Muller



Lectures Outline :

Cloud fundamentals - global distribution, types, visualization and link with large scale circulation

Cloud Formation and Physics - thermodynamics, cloud formation, instability, life cycle of an individual cloud

Organization of deep convection at mesoscales - MCSs, MCCs, Squall lines, Tropical cyclones, Processes, Self-aggregation

Response of the hydrological cycle to climate change - mean precip, precip extremes

Clouds in a changing climate – climate sensitivity, cloud effect, cloud feedback, FAT

Cloud formation



Courtesy : Octave Tessiot

Atmospheric thermodynamics: instability

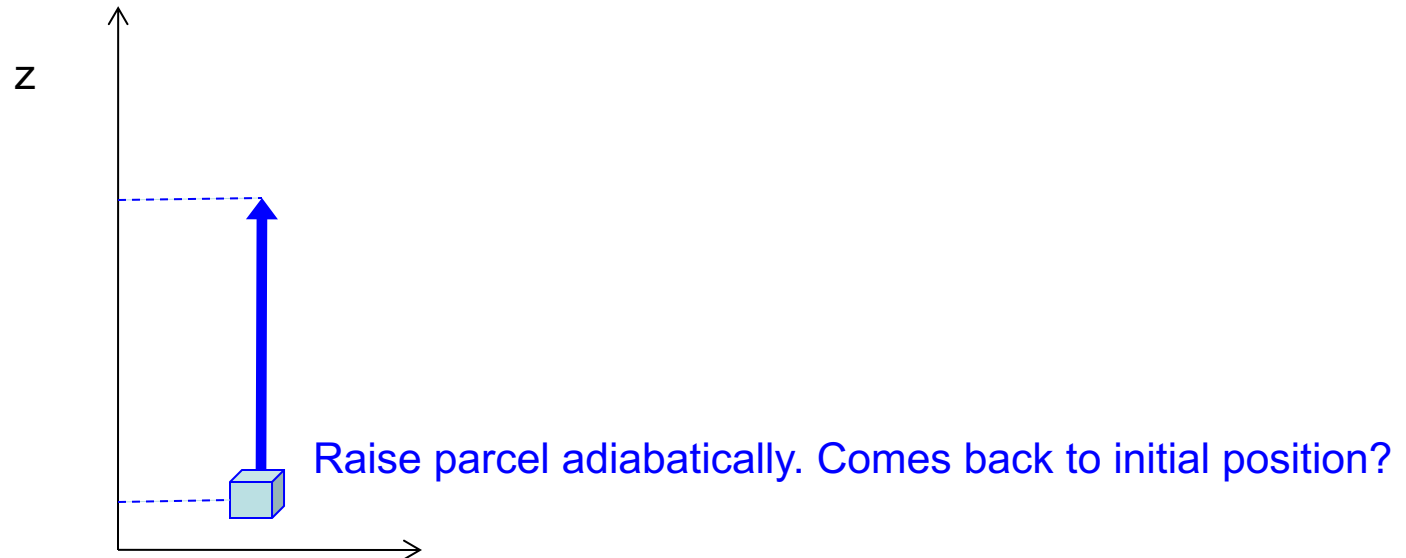
Dry convection

T decreases with height.

But p as well.

Density = $\rho(T,p)$.

How determine stability? The parcel method



Atmospheric thermodynamics: instability

Dry convection

Potential temperature $\theta = T (p_0 / p)^{R/c_p}$ conserved under adiabatic displacements :

Adiabatic displacement

1st law thermodynamics: $d(\text{internal energy}) = \Delta Q$ (heat added) $- \Delta W$ (work done by parcel)

$$c_v dT = - p d(1/\rho)$$

$$\text{Since } p = \rho R T, \quad c_v dT = - p d(R T / p) = - R dT + R T dp / p$$

$$\text{Since } c_v + R = c_p, \quad c_p dT / T = R dp / p$$

$$\Rightarrow d \ln T - R / c_p d \ln p = d \ln (T / p^{R/c_p}) = 0$$

$$\Rightarrow T / p^{R/c_p} = \text{constant}$$

$\Rightarrow \theta = T (p_0 / p)^{R/c_p}$ **potential temperature is conserved** under adiabatic (reversible) displacement

Remark1: ideal gas law: $pV = Nkt \Leftrightarrow p = \rho R T$, $R=k/m$ where m =molecular mass

Remark2: $c_p = c_v + R > c_v$

Remark3: We assumed $p_{\text{parcel}} = p_{\text{environment}} \Leftrightarrow$ quasistatic displacement

Remark4: If we make the **hydrostatic** approximation, **dry static energy $h = c_p T + g z$ is conserved** :
 $c_p dT / T = R dp / p \Leftrightarrow c_p dT = R T dp / p = - g dz \Leftrightarrow c_p T + g z = \text{constant}$

Atmospheric thermodynamics: instability

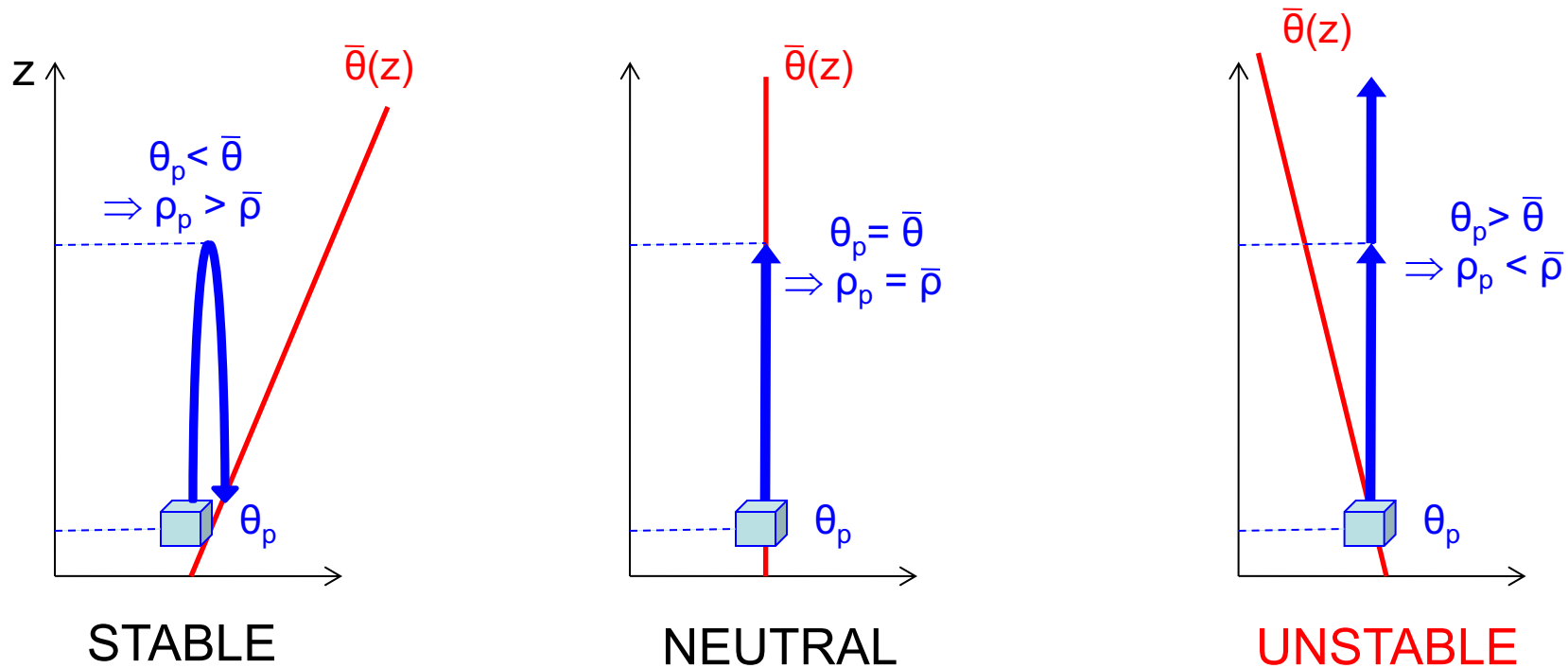
When is an atmosphere unstable to dry convection?

When potential temperature $\theta = T (p_0 / p)^{R/c_p}$ decreases with height !

The parcel method:

Small vertical displacement of a fluid parcel adiabatic ($\Rightarrow \theta = \text{constant}$).

During movement, pressure of parcel = pressure of environment.

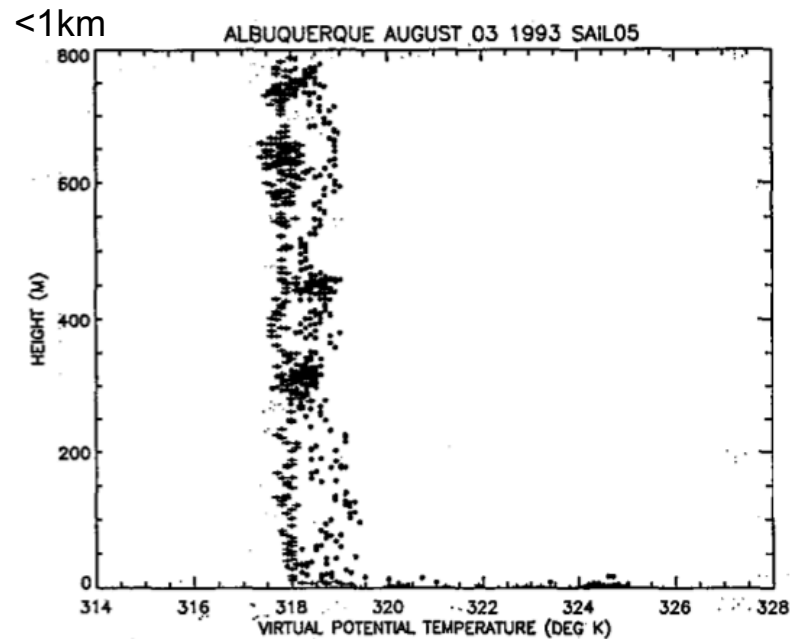


Atmospheric thermodynamics: instability

Convective adjustment time scales is very fast (minutes for dry convection) compared to destabilizing factors (surface warming, atmospheric radiative cooling...)

=> The **observed state is very close to convective neutrality**

Dry convective boundary layer over daytime desert



[Renno and Williams, 1995]

Atmospheric thermodynamics: instability

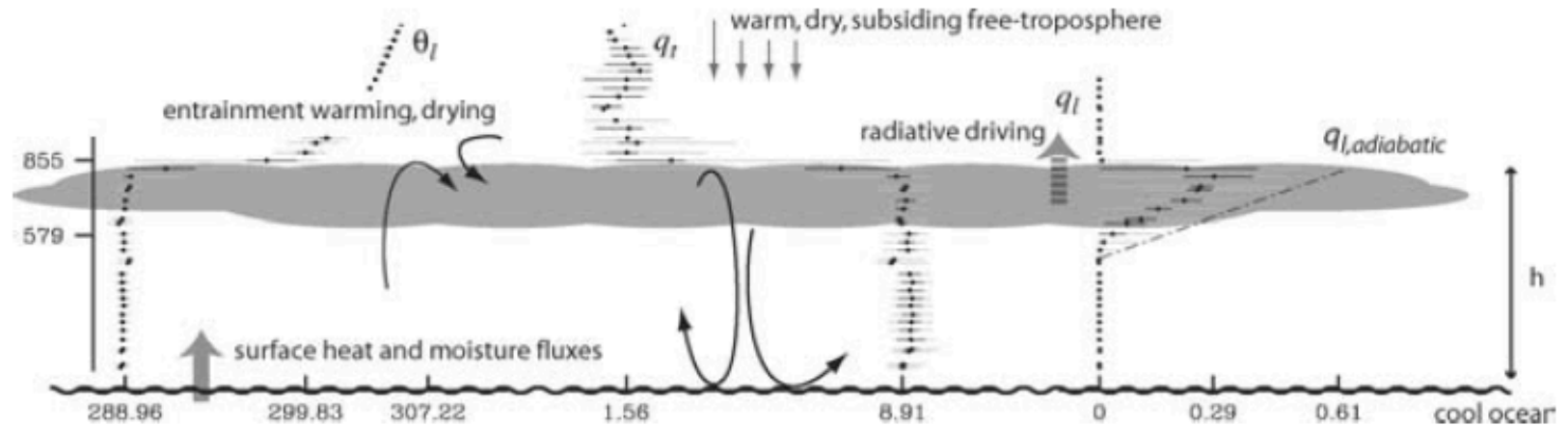


Figure 4 Cartoon of well-mixed, nonprecipitating, stratocumulus layer, overlaid with data from research flight 1 of DYCOMS-II. Plotted are the full range, middle quartile, and mean of θ_l , q_l , and q_l from all the data over the target region binned in 30-m intervals. Heights of cloud base and top are indicated, as are mixed layer values and values just above the top of the boundary layer of various thermodynamic quantities. The adiabatic liquid water content is indicated by the dash-dot line.

Atmospheric thermodynamics: instability

inversion

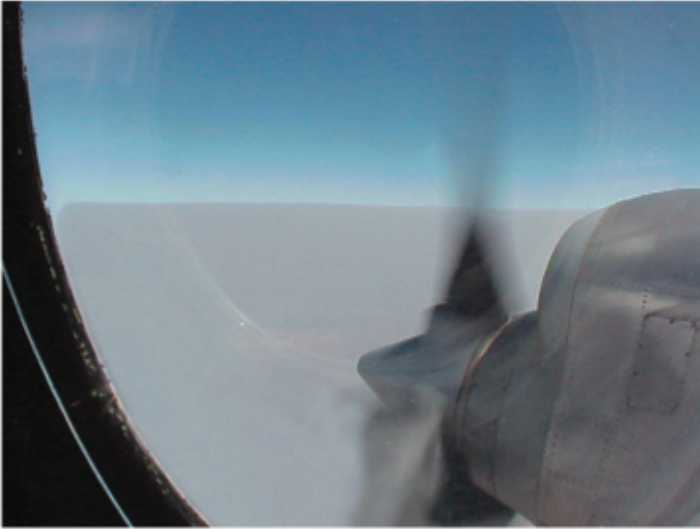


Fig. 3.15 Looking down onto widespread haze over southern Africa during the biomass-burning season. The haze is confined below a temperature inversion. Above the inversion, the air is remarkably clean and the visibility is excellent. (Photo: P. V. Hobbs.)



Smoke rising in [Lochcarron, Scotland](#), is stopped by an overlying layer of warmer air (2006).

Atmospheric thermodynamics: instability

Dry adiabatic lapse rate

We saw that on a dry adiabat, the potential temperature $\theta = T (p_0 / p)^{R/c_p}$ is constant. If in addition we make the hydrostatic approximation, we can deduce the dry adiabatic lapse rate $\Gamma_d = - dT/dz$

Recall in that case $c_p T + g z = \text{constant}$

$$\Leftrightarrow dT/dz = - g / c_p$$

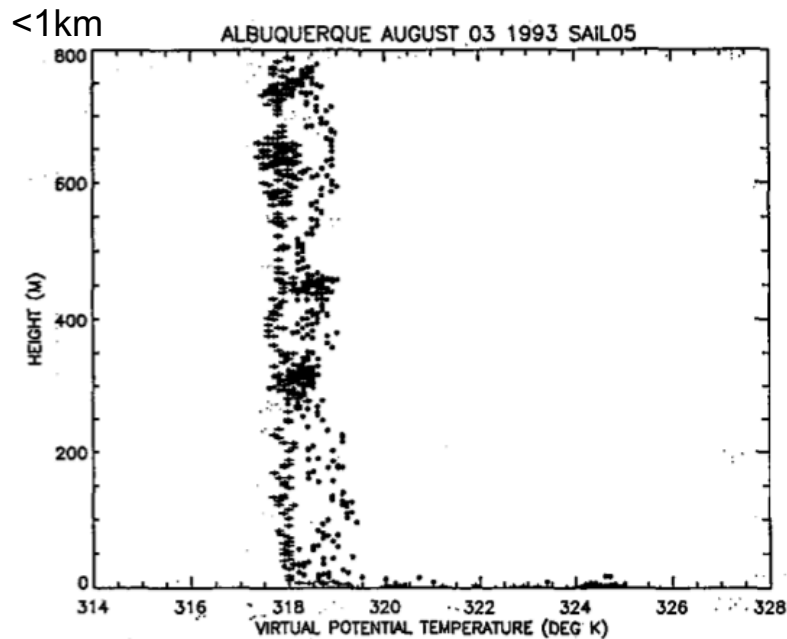
$$\Leftrightarrow \Gamma_d = g / c_p$$

Atmospheric thermodynamics: instability

Convective adjustment time scales is very fast (minutes for dry convection) compared to destabilizing factors (surface warming, atmospheric radiative cooling...)

=> The **observed state is very close to convective neutrality**

Dry convective boundary layer over daytime desert



[Renno and Williams, 1995]

But above a thin boundary layer, not true anymore that $\theta = \text{constant}$. Why?...

Most atmospheric convection involves phase change of water

Significant latent heat with phase changes of water = **Moist Convection**

Atmospheric thermodynamics: instability

Moist variables:

ρ_v = M_v/V water vapor density

ρ_d = M_d/V dry air density

ρ = $\rho_v + \rho_d$ total air density

q_v = ρ_v / ρ water vapor specific humidity

r = ρ_v / ρ_d water vapor mixing ratio

e = partial pressure of water vapor = $\rho_v R_v T$ (ideal gas law for water vapor)

p_d = partial pressure of dry air = $\rho_d R_d T$ (ideal gas law for dry air)

p = $p_d + e$ total pressure (Dalton's law)

T_d = dew point temperature : T at which a parcel must be cooled at constant pressure to reach saturation

T_v = virtual temperature : T that dry air would have to have the same density as moist air at same pressure

Question 1 : Is moist air lighter or heavier than dry air ? In other words is T_v greater or smaller than T ?

Question 2 : Express T_v as a function of T , q_v and $R_d/R_v = \epsilon$

Atmospheric thermodynamics: instability

Question 1 : *Is moist air lighter or heavier than dry air ? In other words is T_v greater or smaller than T ?*

Let's consider a volume of air V at pressure p and temperature T . The ideal gas law implies that $pV=NkT$ where N is the number of molecules in V . So regardless of whether the air is moist or not, the number of molecules is the same. In other words, *moist air is not formed by adding water molecules to the air, but by replacing dry air molecules with water molecules.*

If we compare the molecular masses of H_2O , N_2 and O_2 , clearly H_2O is the lightest :
 $m_{H_2O}=(2+16)m_H$; $m_{O_2}=2*16m_H$; $m_{N_2}=2*14m_H$.

So moist air is lighter than dry air. Hence $T_v > T$.

In fact the ratio of molecular masses is $m_v/m_d \sim .622 = \epsilon$

In atmospheric applications we use the ideal gas law with density :

$$pV=NkT \Leftrightarrow p = (N m / V) (k/m) T = \rho R T$$

Question 2 : *Express T_v as a function of T , q_v and $R_d/R_v = \epsilon$*

By definition, T_v satisfies $p = \rho R_d T_v$ with $p = e + p_d = (\rho_v R_v + \rho_d R_d) T$.

$$\text{So } T_v = \left\{ \frac{\rho_v}{\rho} \frac{R_v}{R_d} + \frac{(\rho - \rho_v)}{\rho} \right\} T = \left\{ \frac{q_v}{\epsilon} + 1 - q_v \right\} T$$

$$\Rightarrow T_v = \left\{ 1 + \left(\frac{1}{\epsilon} - 1 \right) q_v \right\} T \sim (1 + .61 q_v) T > T \text{ as expected.}$$

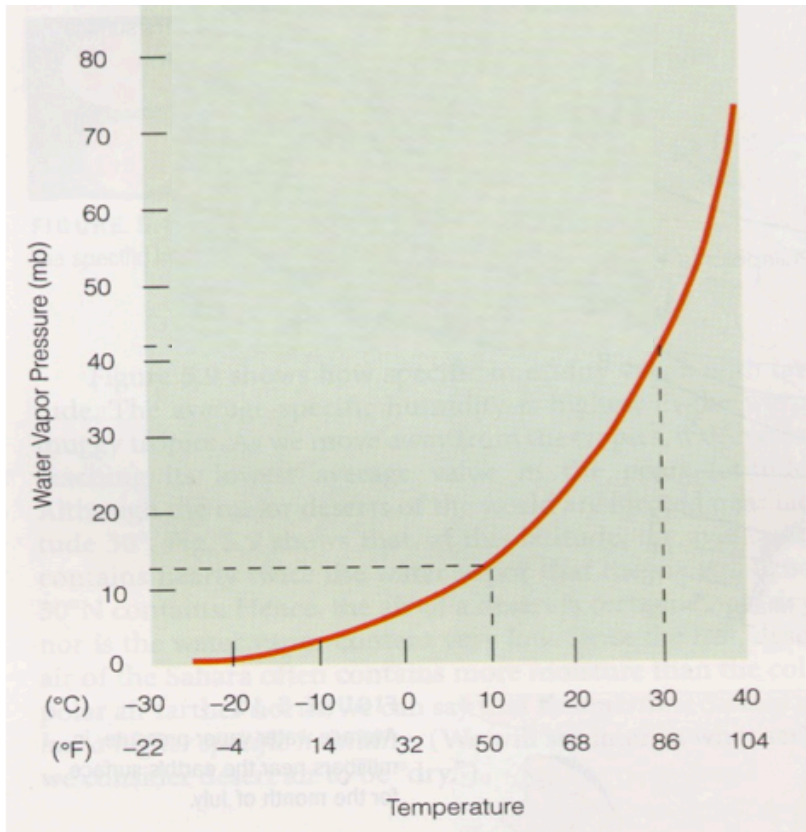
Atmospheric thermodynamics: instability

Clausius Clapeyron
$$\frac{de_s}{dT} = \frac{L_v(T)e_s}{R_v T^2}$$

$$e_s(T)$$

where:

- e_s is saturation vapor pressure,
- T is a temperature,
- L_v is the specific latent heat of evaporation,
- R_v is water vapor gas constant.



e_s depends only on temperature

e_s increases roughly exponentially with T

Saturation water vapor amount increases with temperature

RH = relative humidity = e/e_s

Remark on « sponge theory » :

« Warm air can hold more water vapor than cold air » ...

Has nothing to do with air, similar in other gas !

CC \Leftrightarrow Equilibrium between condensation and evaporation

Atmospheric thermodynamics: instability

When is an atmosphere unstable to moist convection ?

Equivalent potential temperature $\theta_e = T (p_0 / p)^{R/c_p} e^{L_v q_v / (c_p T)}$ is *approximately* conserved under adiabatic displacements :

1st law thermodynamics if air saturated ($q_v=q_s$) :

$$d(\text{internal energy}) = Q (\text{latent heat}) - W (\text{work done by parcel})$$

$$c_v dT = - L_v dq_s - p d(1/\rho)$$

$$\Rightarrow d \ln T - R / c_p d \ln p = d \ln (T / p^{R/c_p}) = - L_v / (c_p T) dq_s$$

$$\Rightarrow T / p^{R/c_p} e^{L_v q_s / (c_p T)} \sim \text{constant}$$

Note: Air saturated $\Rightarrow q_v=q_s$
Air unsaturated $\Rightarrow q_v$ conserved

Hence

$\theta_e = T (p_0 / p)^{R/c_p} e^{L_v q_v / (c_p T)}$ equivalent potential temperature is conserved

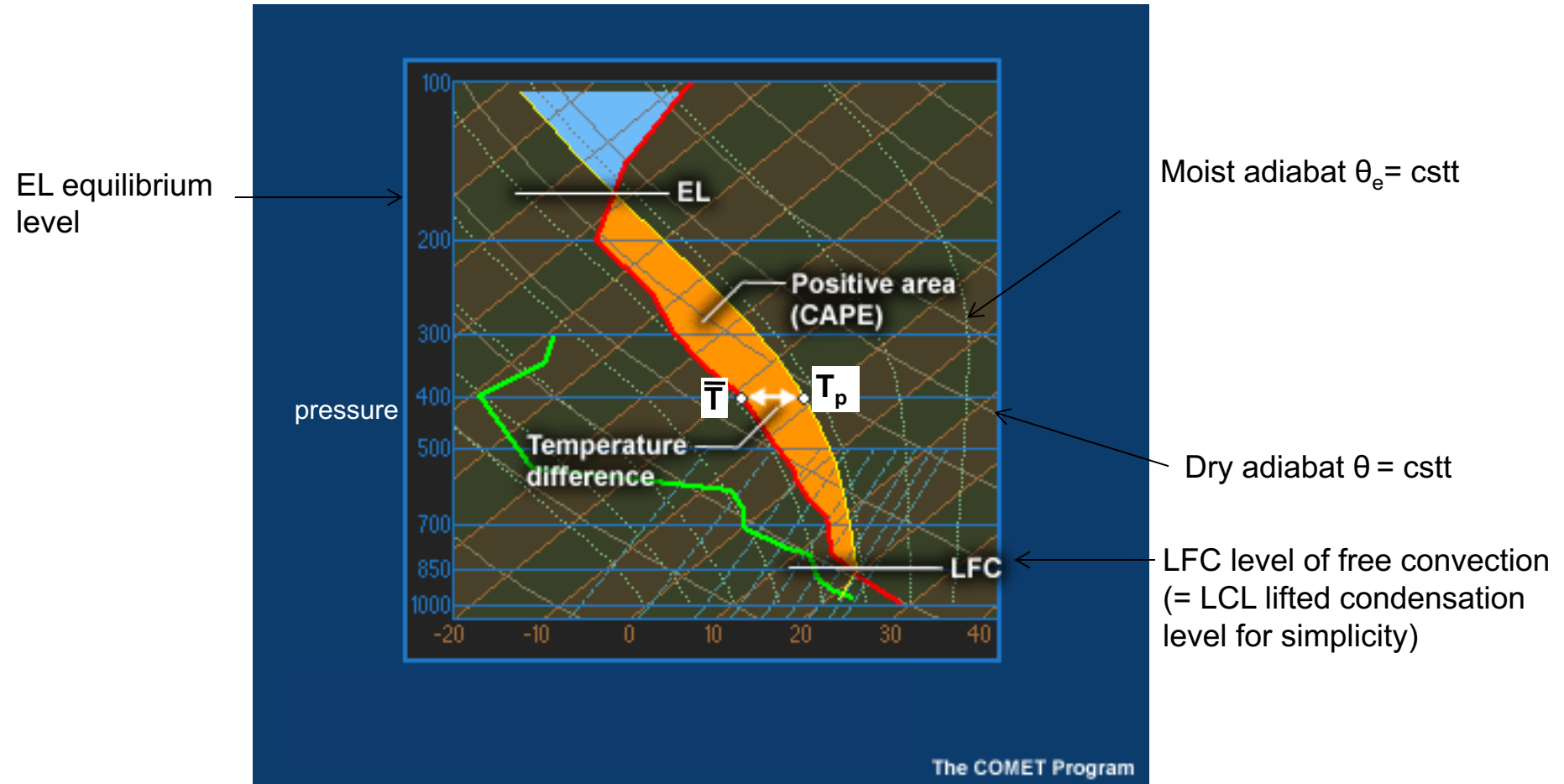
Remark : If we make the **hydrostatic** approximation, **MOIST STATIC ENERGY** $h = c_p T + g z + L_v q_v$ is conserved :

$$c_p dT - R T dp / p = - L_v dq_s \Leftrightarrow c_p dT + g dz = - L_v dq_s \Leftrightarrow c_p T + g z + L_v q_v = \text{constant}$$

Atmospheric thermodynamics: instability

When is an atmosphere unstable to moist convection ?

Skew T diagram (isoT slanted), atmospheric T in red

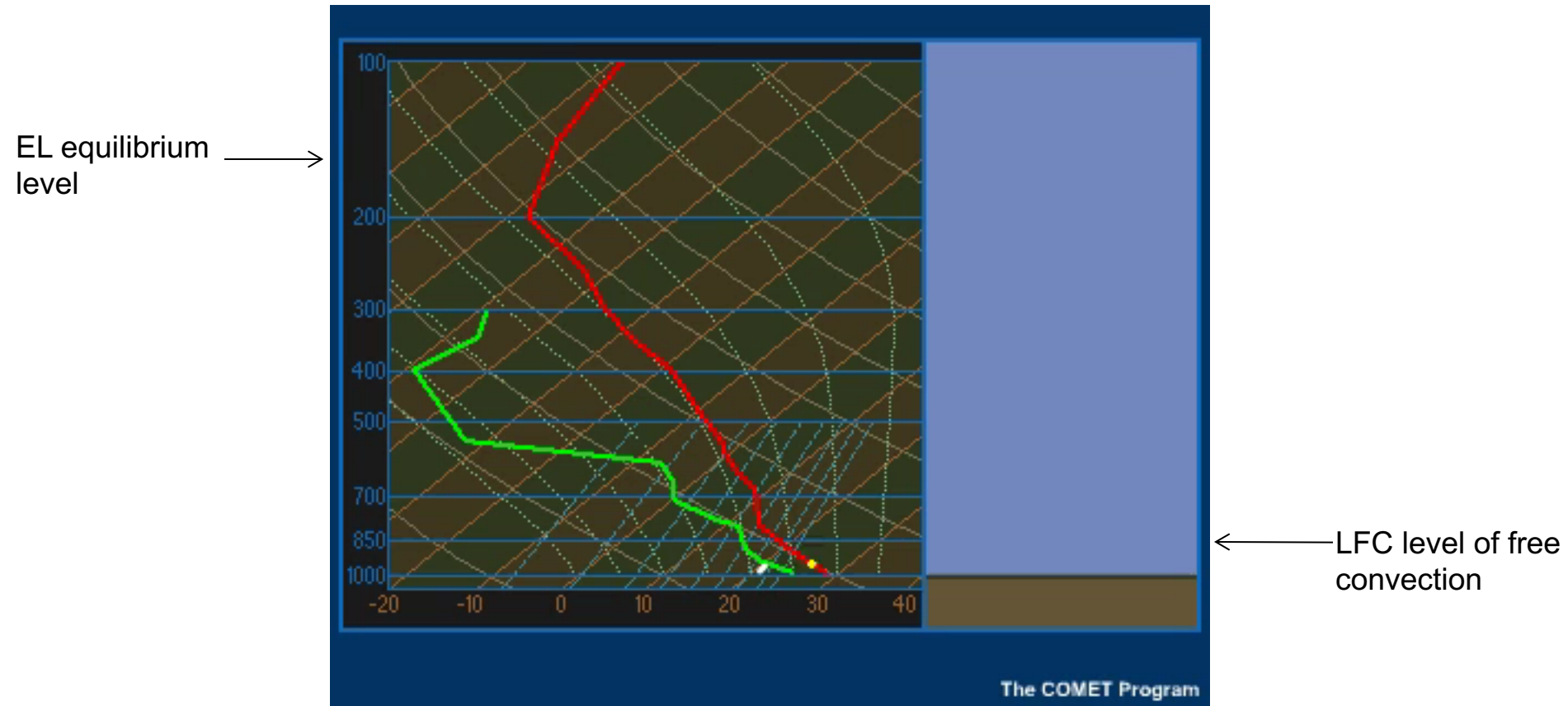


CAPE: convective available potential energy

Atmospheric thermodynamics: instability

Moist convection

Parcel = yellow dot

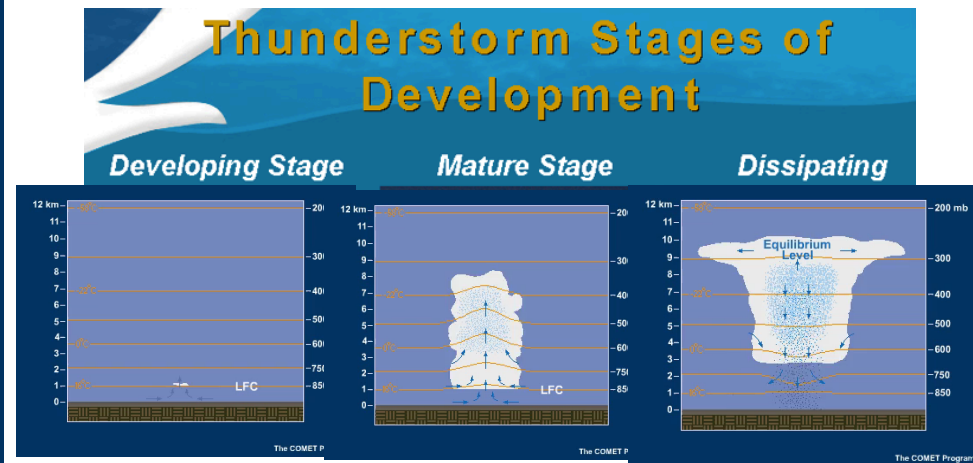
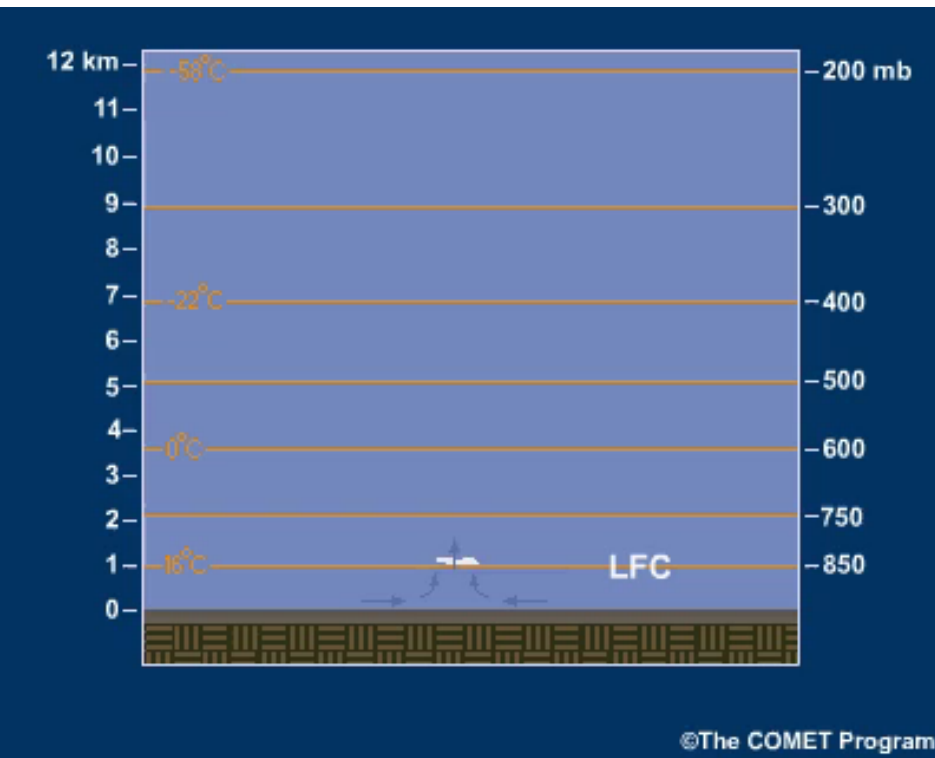


CAPE: convective available potential energy

Atmospheric thermodynamics: instability

If enough atmospheric instability present, cumulus clouds are capable of producing serious storms!!!

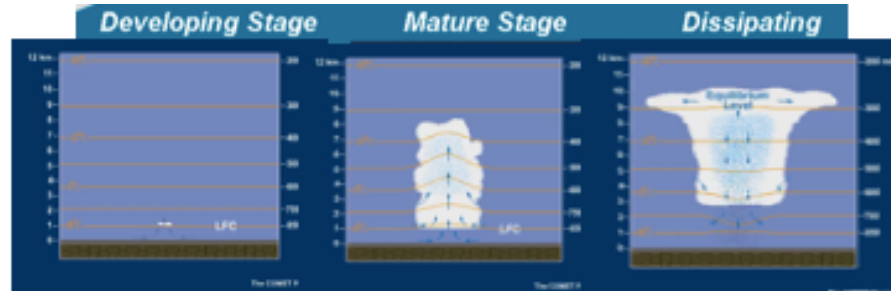
Strong updrafts develop in the cumulus cloud => mature, deep cumulonimbus cloud. Associated with heavy rain, lightning and thunder.



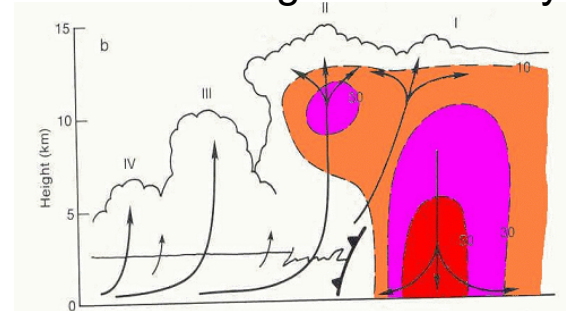
↗
Evaporative driven cold pools

Atmospheric thermodynamics: instability

Note that thunderstorms can be : single-cell (typically with weak wind shear)

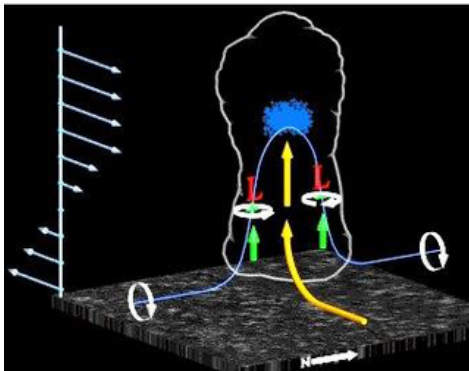
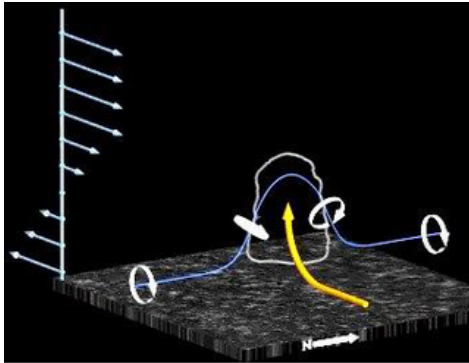
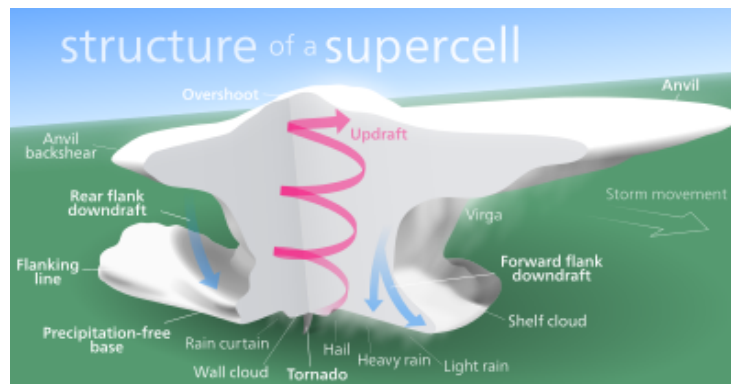


multi-cell (composed of multiple cells, each being at a different stage in the life cycle of a thunderstorm).



or supercell, characterized by the presence of a deep, rotating updraft

Typically occur in a significant vertically-sheared environment



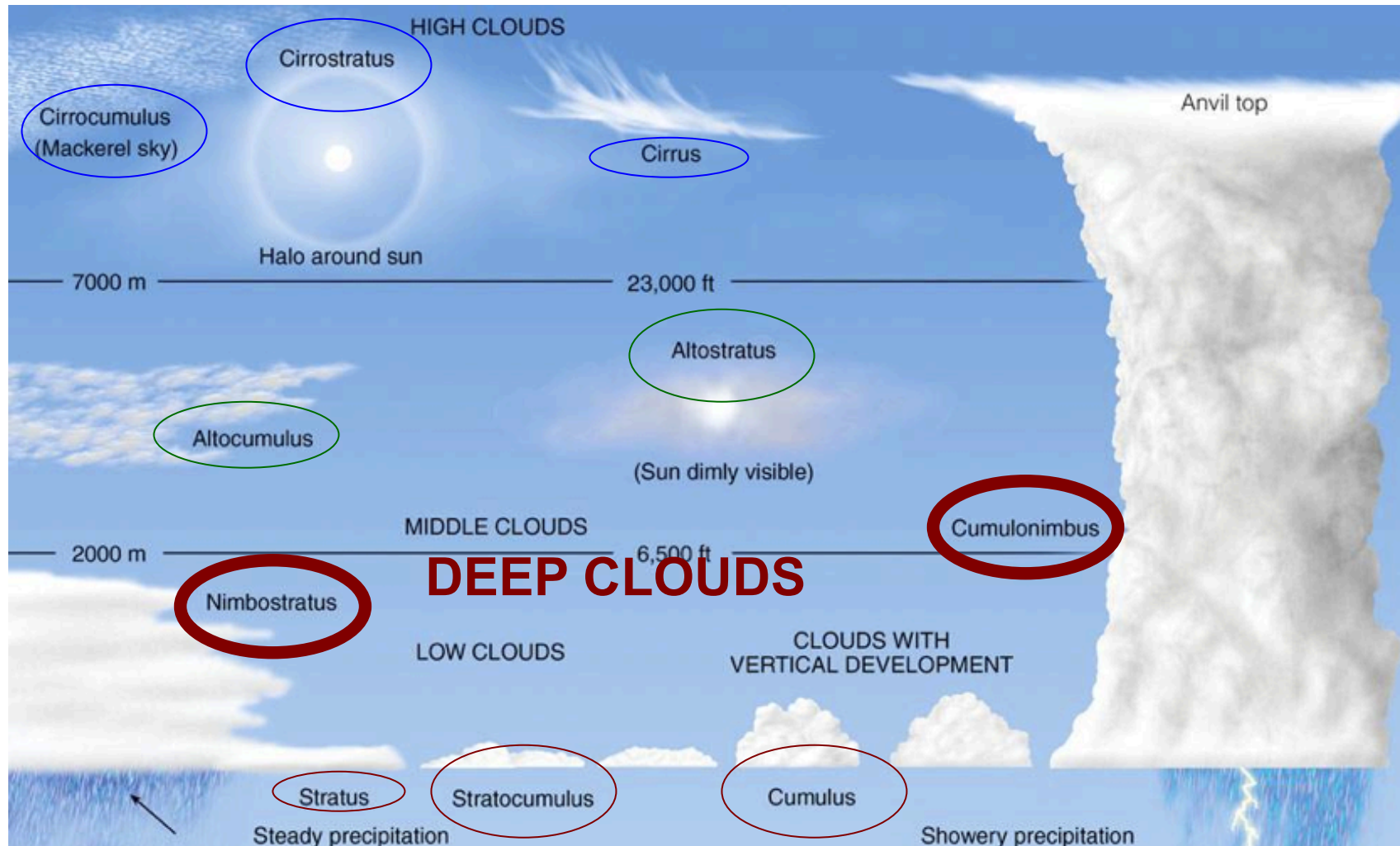
Processes leading to cloud formation

How do those physical considerations explain cloud formation ?

=> FOR DEEP CLOUDS :

We saw that for deep clouds, adiabatic ascent from an unstable BL parcel (warm and/or moist) *rising through an unstable atmospheric T profile* can lead to strong deep convection.

Other lifting mechanisms ?



Cloud formation: Deep clouds

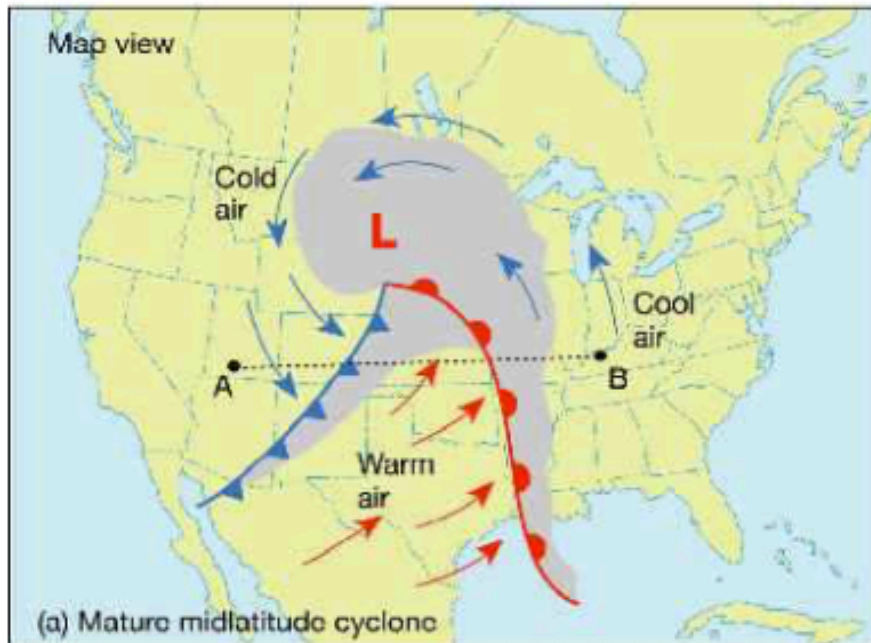
Other lifting mechanisms :

- orography
- large-scale convergence
- fronts

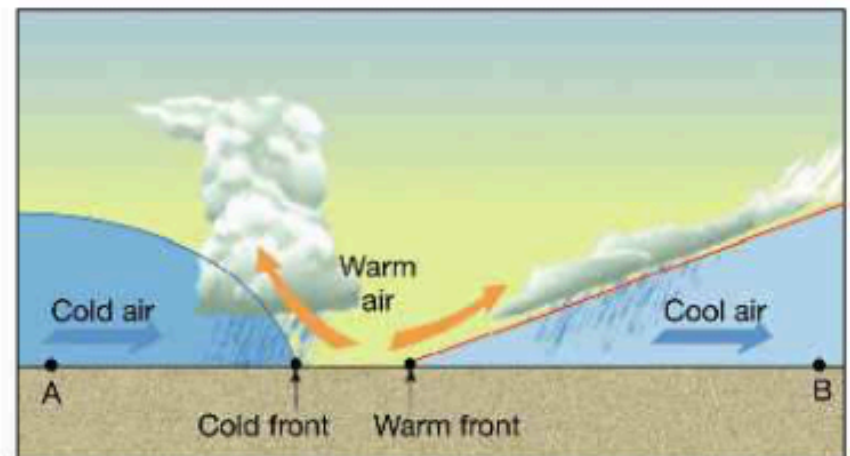
=> All force ascent, and leads to *deep convection if atmosphere above is unstable*

Clouds associated with a frontal system

(blue : cold front, steep and fast; red: warm front, shallower and slower)



Cross sectional view

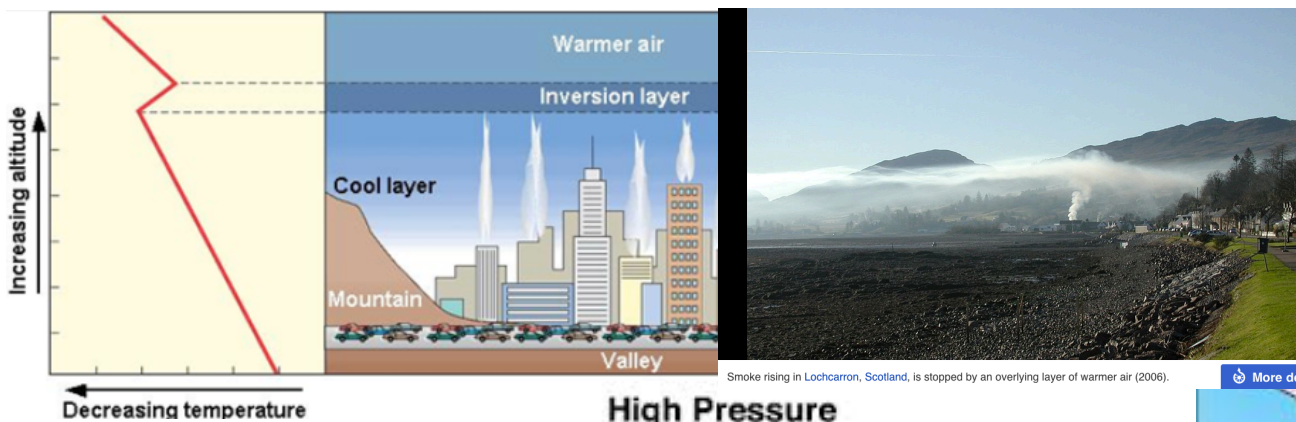


Cloud formation: Shallow layer clouds

SHALLOW LAYER CLOUDS

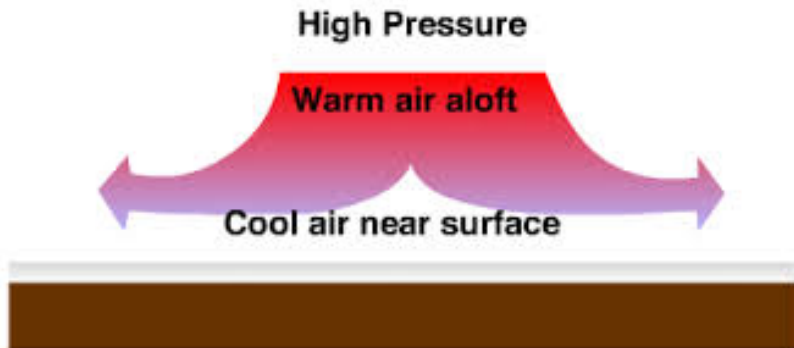
- **Fog and stratus:** in BL cooled from below, by radiation or conduction from cold surface => **Stable BL**, reach saturation by cooling
- **Stratus or stratocumulus or shallow cumulus:** in BL heated from below => **Unstable BL, with a stable atmosphere above**. Also radiative cooling at the top of the cloud layer destabilizes the layer and contributes to the convection.

When do we have unstable layer capped by stable layer ? **Warm** air above **cool** air « T inversion »
An inversion can develop aloft as a result of air gradually sinking over a wide area and being warmed by adiabatic compression, e.g. associated with subtropical high-pressure areas.



Smoke rising in Lochcarron, Scotland, is stopped by an overlying layer of warmer air (2006).

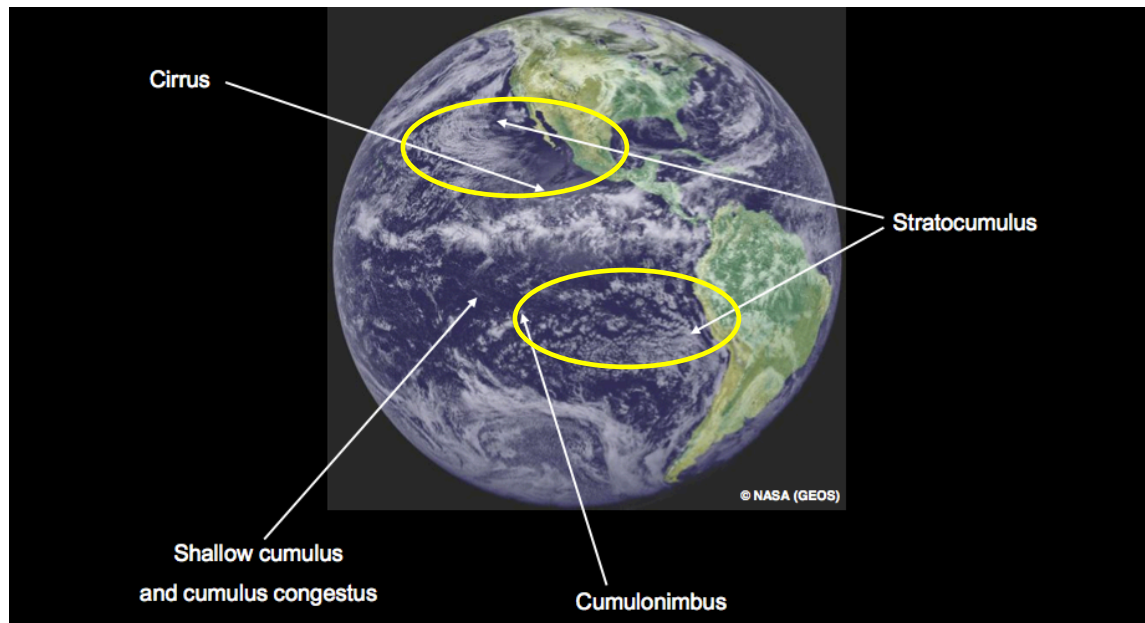
[More det](#)



Cloud formation: Shallow layer clouds

SHALLOW LAYER CLOUDS

- **Fog and stratus**: in BL cooled from below, by radiation or conduction from cold surface
=> **Stable BL**, reach saturation by cooling
- **Stratus or stratocumulus or shallow cumulus**: in BL heated from below
=> **Unstable BL, with a stable atmosphere above**. Also radiative cooling at the top of the cloud layer destabilizes the layer and contributes to the convection.
e.g. - subtropical latitudes **west of continents** stratus and stratocumulus associated with anticyclones around high pressure
 - middle and high latitudes cold air offshore across the coastlines of cold continents or ice sheets, over warm ocean => stratocumulus

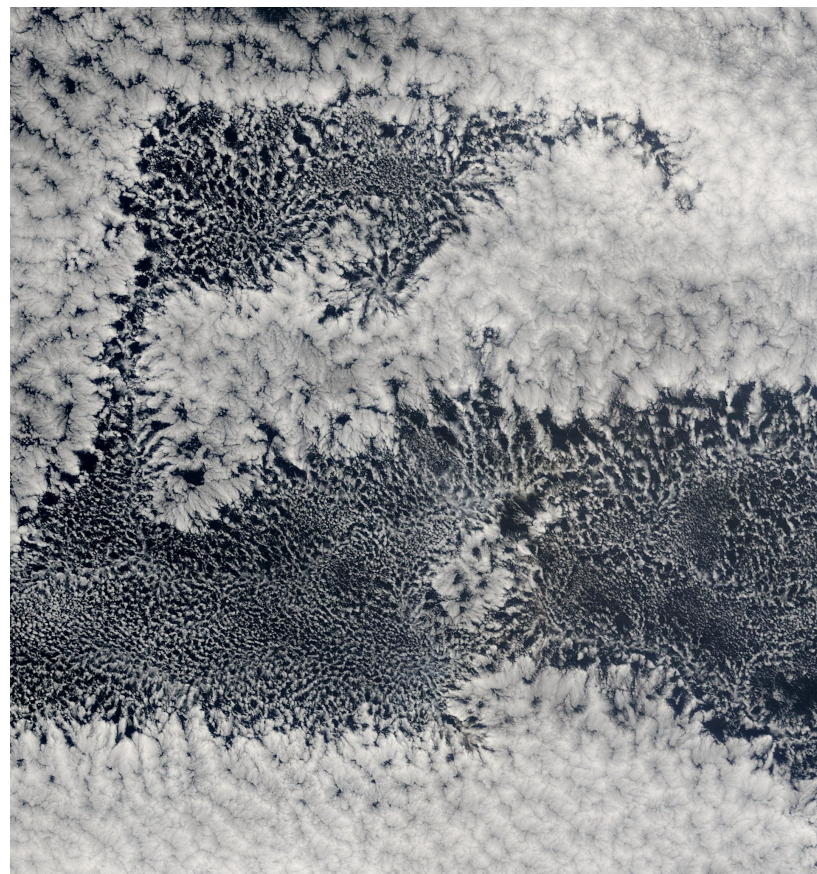
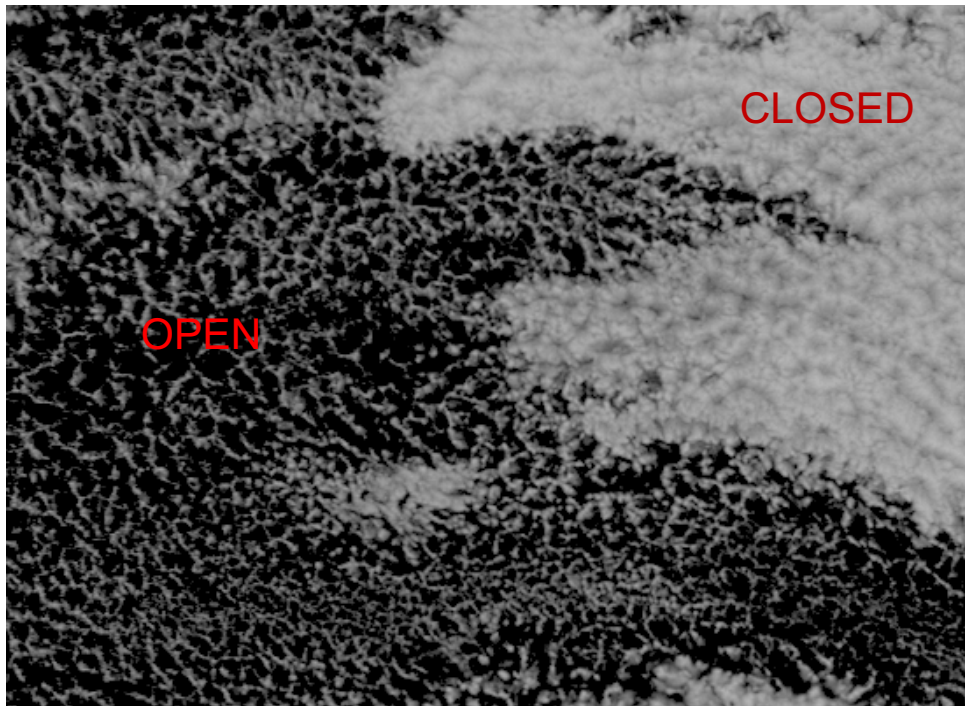


Cloud formation: Shallow layer clouds

MESOSCALE ORGANIZATION OF THOSE CLOUDS : Open and closed cells

Not all is known, but some processes appear to play a role : shear, thermal instabilities of the BL, cloud-top entrainment and precipitation-driven cold pools (aerosols ?)

Mesoscale cellularity in marine stratocumulus clouds. This MODIS image (approximately 800 km across) shows the sharp transitions that occur between the closed and open cells.



Cloud formation: Shallow layer clouds

Mesoscale organization of shallow clouds : Open and closed cells shallow convection

172

5 Shallow-Layer Clouds

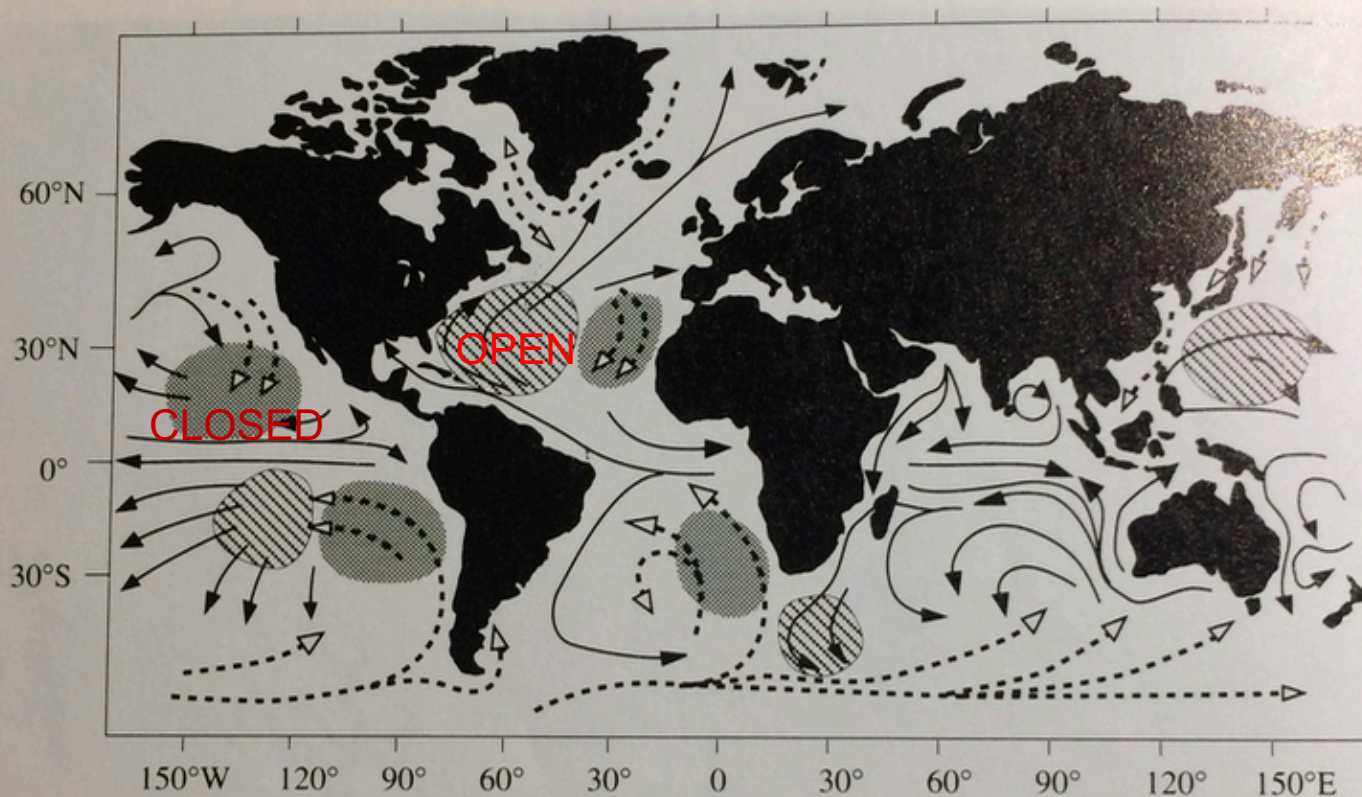
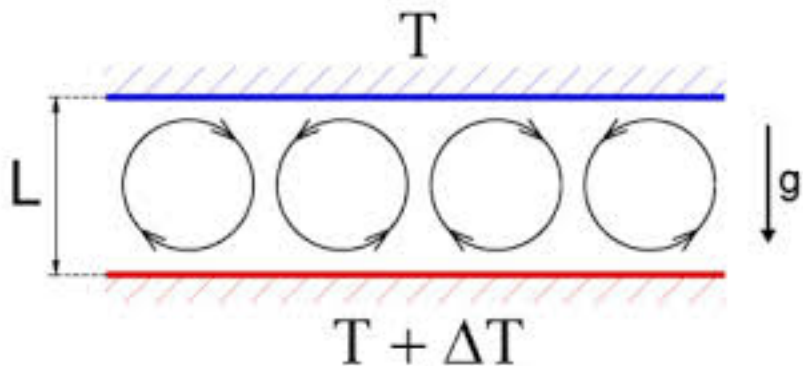
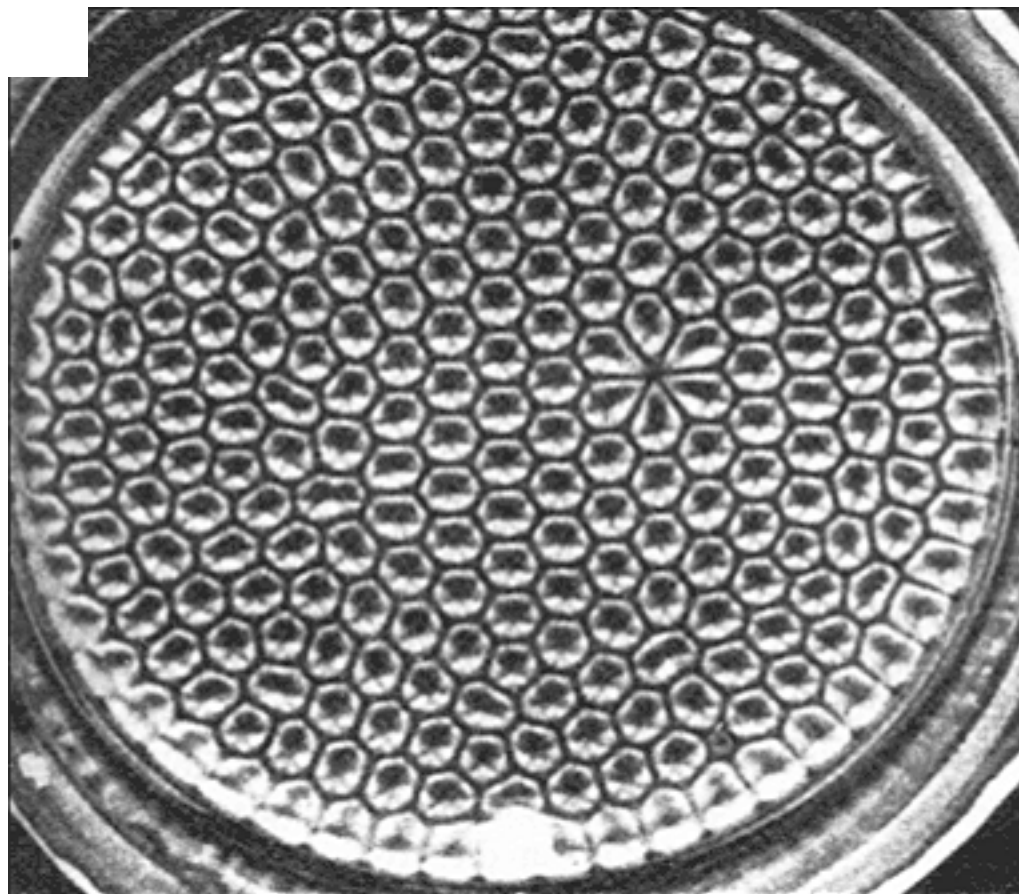


Figure 5.24 Global climatology of cellular structure of stratocumulus and small cumulus over oceans. Shaded areas are regions where closed cells predominate. Hatched areas show where open cells are more common. Solid streamlines show locations of warm ocean currents. Dashed streamlines show cold currents. Land masses are blackened. (Adapted from Agee *et al.*, 1973. Reproduced with permission from the American Meteorological Society.)

Cloud formation: Shallow layer clouds



Rayleigh-Benard convection



Cloud formation: Shallow layer clouds

Rayleigh-Benard convection



Gold paint dissolved in acetone. Put it in a shallow dish. Cover it so that the acetone does not evaporate. Let stabilise and then remove the cover. Evaporation of the acetone causes the top layer to cool thus starting convection.

Cloud formation: Shallow layer clouds

SHALLOW LAYER CLOUDS

- Cirriform clouds:

Not much water vapor at those high altitudes => mainly radiation driven.

Clouds of (mainly) ice in an **unstable layer between two stable layers**

SW heating throughout the clouds, while LW cools above and warms below

-Can be detrained from deep convective clouds (most often, consistent with largest cirrus cover in the tropics and in the extratropics where deep convection), or

-Can occur away from generating source when unstable layer aloft

- Altostratus & altocumulus: these can be

-Remnants of other clouds: protruding layers in middle levels due to horizontal wind

-Altocumulus also sometimes high-based convective clouds => same dynamics as deep convective clouds

-Altostratus or shallow layer of altocumulus can also resemble a radiatively driven « mixed layer » aloft, leading to a cloud-filled layer radiatively driven at its top
(Can lead to rolls in the absence of shear)

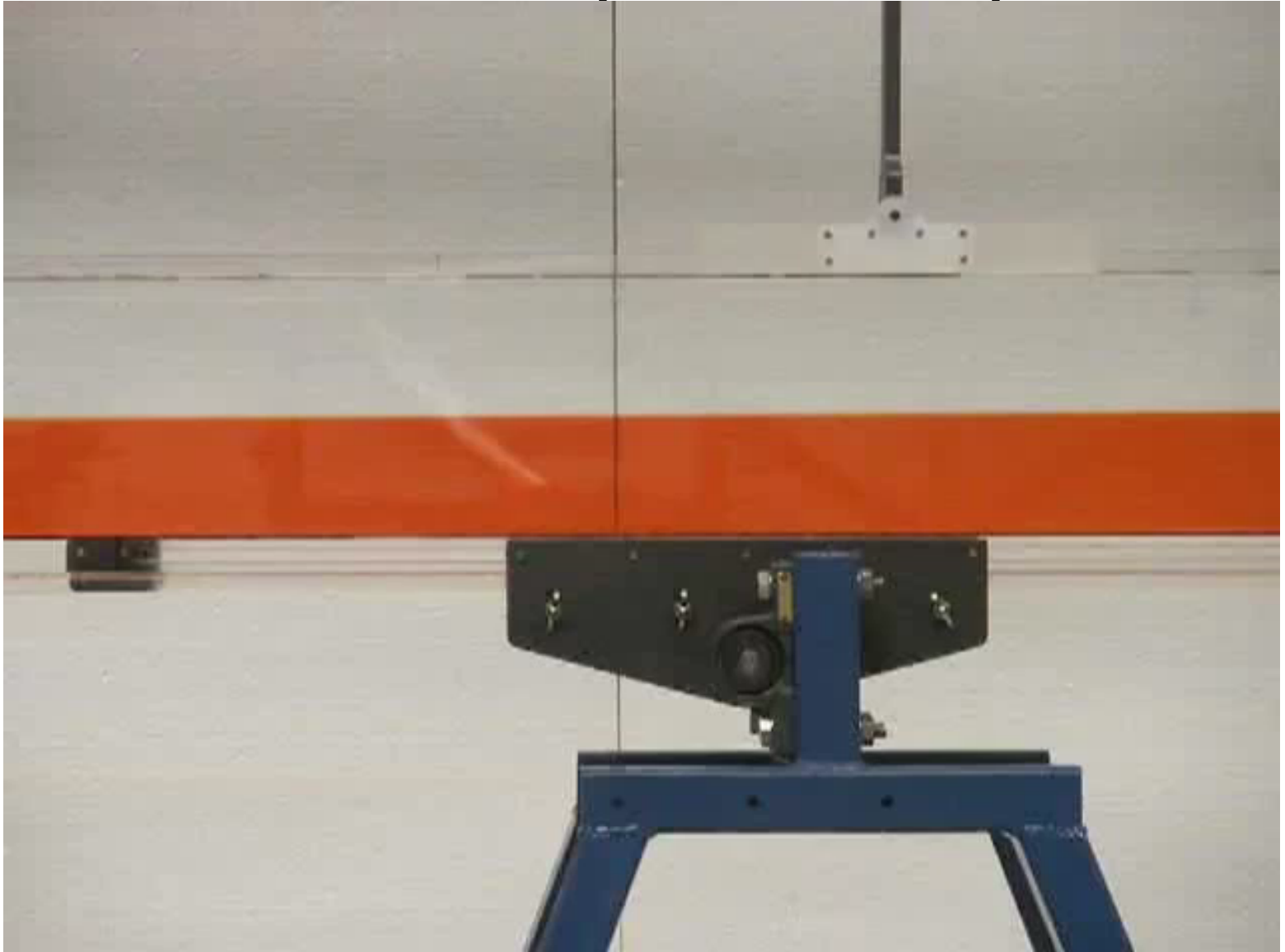
Atmospheric thermodynamics: instability

Kelvin Helmholtz Instability : destabilized by shear



Atmospheric thermodynamics: instability

Kelvin Helmholtz Instability : destabilized by shear



Lectures Outline :

Cloud fundamentals - global distribution, types, visualization and link with large scale circulation

Cloud Formation and Physics - thermodynamics, cloud formation, instability, life cycle of an individual cloud

Organization of deep convection at mesoscales - MCSs, MCCs, Squall lines, Tropical cyclones, Processes, Self-aggregation

Response of the hydrological cycle to climate change - mean precip, precip extremes

Clouds in a changing climate – climate sensitivity, cloud effect, cloud feedback, FAT